## ON THE CONSERVATION OF VORTICITY IN THE MAGNETOHYDRODYNAMICS OF PLASMA WITH MANY COMPONENTS

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The system of equations for the hydrodynamic approximation of an ideal plasma consisting of N kinds of ions (k = 1, ..., N) takes the form [1-4]

$$\frac{\partial \mathbf{v}_{k}}{\partial t} + \nabla \left( \frac{\boldsymbol{v}_{k}^{2}}{2} + \int \frac{d\boldsymbol{p}_{k}}{n_{k}m_{k}} + \boldsymbol{F}_{k} \right) = \mathbf{v}_{k} \times \left( \operatorname{rot} \mathbf{v}_{k} + \frac{\mu \boldsymbol{e}_{k}}{cm_{k}} \mathbf{H} \right) + \frac{\boldsymbol{e}_{k}}{m_{k}} \mathbf{E}$$

$$\frac{\partial n_{k}}{\partial t} + \operatorname{div} n_{k} \mathbf{v}_{k} = 0, \quad \operatorname{div} \mathbf{H} = 0, \quad \operatorname{div} \mathbf{E} = \frac{4\pi}{\varepsilon} \sum_{f=1}^{N} \boldsymbol{e}_{f} n_{f} \quad (1)$$

$$\operatorname{rot} \mathbf{H} = \frac{4\pi}{c} \sum_{f=1}^{N} \boldsymbol{e}_{f} n_{f} \boldsymbol{v}_{f} + \frac{\varepsilon}{c} \frac{\partial \mathbf{E}}{\partial t}, \quad \operatorname{rot} \mathbf{E} = -\frac{\mu}{c} \frac{\partial \mathbf{H}}{\partial t}$$

The first two equations are the momentum and continuity equations for the ions of species k, and the rest are Maxwell's equations. The inertial term in the momentum equation is given in the form of I.S. Gromeka; H and E denote the intensity of the magnetic and electric fields,  $\mu$  the magnetic permeability,  $\epsilon$  the dielectric constant and c the speed of light.  $F_k$  represents the potential of external forces,  $V_k$  the velocity,  $P_k$  the partial pressure,  $n_k$  the number of particles per unit volume,  $e_k$ and  $\mathbf{m}_k$  the charge and the mass, all pertaining to species k.

The plasma is considered either as a "cold"  $(p_k = 0)$  or an incompressible  $(n_k = \text{const})$  or, again, a barotropic  $(p_k = p_k(n_k))$  medium.

When the curl operator is applied to the momentum equation and use is made of the equation for the electric induction, there results

$$\frac{\partial \mathbf{\Omega}_k}{\partial l} = \operatorname{rot} \mathbf{v}_k \times \mathbf{\Omega}_k, \qquad \mathbf{\Omega}_k \equiv \operatorname{rot} \mathbf{v}_k + \frac{\mu e_k}{cm_k} \mathbf{H}$$
(2)

The conservation of the partial vorticity  $\Omega_k$  is analogous to the conservation of vorticity in ordinary hydrodynamics as shown by Helmholtz [4]. Hence, the vortex lines corresponding to  $\Omega_k$  are "glued" to the particles of the k species, and the flux of the partial vorticity  $\Omega_k$ across an arbitrary material surface, made up of particles of species k, is conserved. Furthermore, the circulation of the vector

$$\mathbf{v}_k + (\mu e_k / cm_k) \mathbf{A} (\mathbf{H} = \operatorname{rot} A)$$

along arbitrary closed contours, made up of particles of species k, is also conserved. We remark that the conservation of the partial vorticity ceases as a result of viscous forces and collisions between particles of different species.

For the cold electron-ion plasma the vorticity-conservation laws were noted in [1]. In [2], where the plasma was described by a system of magnetohydrodynamic equations for a binary (electron-ion) medium, the conservation of the partial vorticity for the ionic component was brought out.

For the cases of the steady two-parameter and the steady vortical plasma motions the integral form of the conservation of the partial vorticity was noted in [3].

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